3D microwave simulation for spherical tokamaks

Tom Williams¹ trnw500@york.ac.uk

Acknowledgments:

Roddy Vann¹, Martin O' Brien², Vladimir Shevchenko², Simon Freethy², Alf Köhn³

¹York Plasma Institute, Department of Physics, University of York, Heslington, York YO10 5DD, UK ²EURATOM/CCFE Fusion Association, Culham Science Centre, Abingdon, Oxon OX14 3DB, UK ³IGVP, Universität Stuttgart, Pfaffenwaldring 31, D-70569 Stuttgart, Germany











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- 1. Why study microwave interactions?
- 2. Underlying plasma physics
- 3. 3D full-wave simulations
- 4. Ongoing work



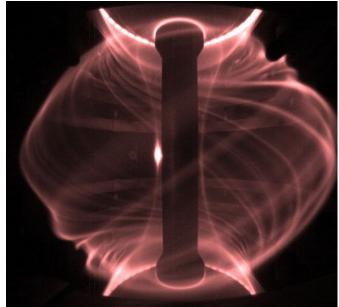
3D effects

• Spherical tokamak edge plasma contains 3D density fluctuations (filaments, blobs etc.) and magnetic shear

• Interactions with microwaves must be understood for EC emission diagnostics, heating and current drive (for EBW, the effect on mode conversion)

• 3D full-wave modelling necessary to explore interactions in detail, investigate new physics and aid interpretation of experimental data

• Extrapolation beyond current experiments

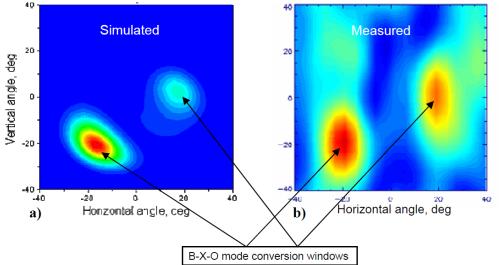


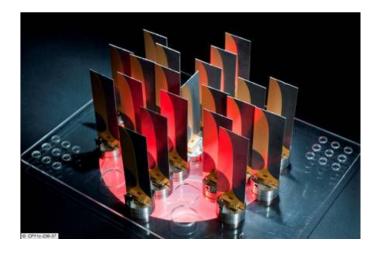
MAST plasma showing filaments at edge



SAMI diagnostic

- York/Culham collaboration (V. Shevchenko et al., 2012, arXiv: 1210.3278 [physics.plasm-ph])
- Images microwave emission at 10 35 GHz \rightarrow radial range through edge. High time resolution (~10 μ s)
- In process of using this data to generate an edge **J**-profile – aim to reconstruct pedestal during inter-ELM period. Major H-mode issue
- Observed fluctuations much higher than expected! (Dave had a poster...)







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Cold plasma dispersion

• Assume plane waves & rewrite linearised Maxwell's equations using dielectric permittivity tensor . Matrix equation obtained:

$$\left[\mathbf{k}\mathbf{k} - k^{2}\underline{I} + \frac{\omega^{2}}{c^{2}}\underline{\varepsilon}\right] \cdot \mathbf{E} = \underline{M} \cdot \mathbf{E} = 0$$

• Ideal, cold, magnetised plasma with uniform equilibrium B_0 -field. Evaluate dielectric tensor using linearised fluid equation for electrons:

$$m_e \frac{\partial \mathbf{v}_e}{\partial t} = -e \left(\mathbf{E} + \mathbf{v}_e \times \mathbf{B}_0 \right)$$

• Assume \mathbf{B}_0 and \mathbf{k} are perpendicular (i.e. propagation \perp background magnetic field) – find 2 solutions to matrix equation:

$$n^{2} = 1 - X$$
 $n^{2} = 1 - \frac{X(1 - X)}{1 - X - Y^{2}}$ $\left(X = \frac{\omega_{pe}^{2}}{\omega^{2}}, Y = \frac{|\omega_{ce}|}{\omega}\right)$

O-mode

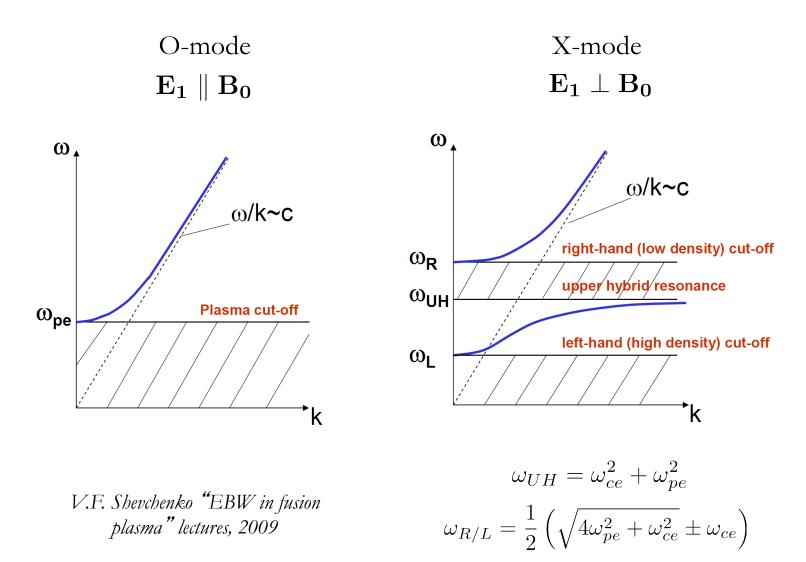
X-mode

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Cold plasma modes



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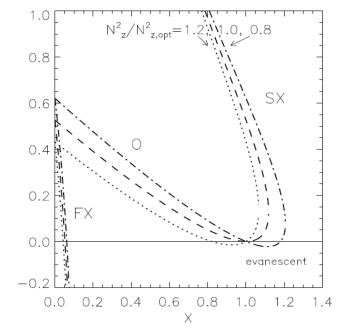
X-O conversion

• Mode conversion from X-mode to O-mode (and vice versa) occurs at O-mode density cut-off ($\omega = \omega_{pe}$) if wave is obliquely incident at an optimal angle to **B**₀

- At suboptimal angles, wave tunnels through evanescent layer of finite width (dependent on density scale length L_n), reducing conversion efficiency
- Using WKB approximation, this efficiency T was calculated by Mjølhus :

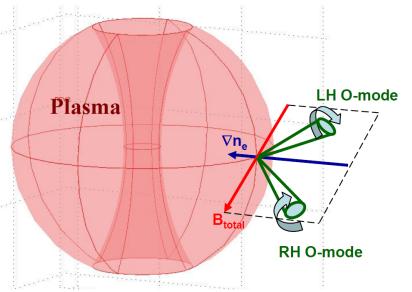
$$T = \exp\left\{-\pi k_0 L_n \sqrt{\frac{Y}{2}} \left[2(1+Y)(N_{z,opt} - N_z)^2 + N_y^2\right]\right\}$$

E. Mjølhus, J. Plasma Physics 31 (1) 7, 1984



B-X-O in a spherical tokamak

- Electrostatic electron Bernstein modes (EBWs) are excited near cyclotron resonances and couple to X-mode at the upper hybrid resonance
- Typically, in STs such as MAST running at high β (higher n_e, lower **B**₀), $\omega_{pe} > \omega_{ce}$
- Problematic for conventional ECE diagnostics but allows B-X-O conversion to produce two cones of O-mode emission fror
- Cones emitted in the plane of $\mathbf{B_0}$ and ∇n_e
- ∇n_e known from TS diagnostic
- Imaging these cones gives pitch of $\mathbf{B}_0 \rightarrow \mathbf{B}_{\theta} \rightarrow \mathbf{J}$ at edge

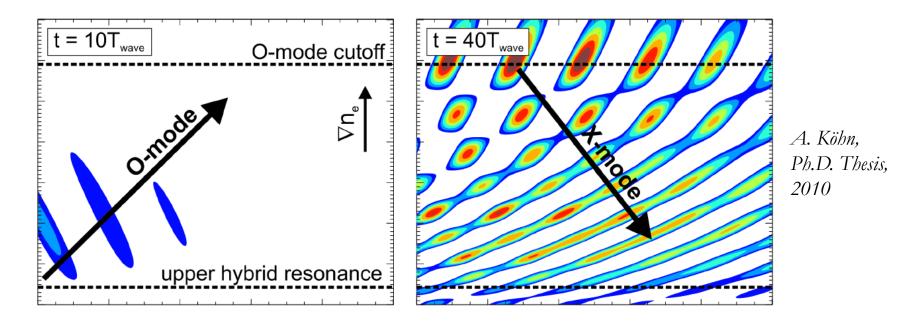




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Previous modelling



- Ray/beam tracing applied to beam propagation problems, but fast variations in refractive index make it unsuitable for conversion region
- 2D full-wave modelling of O-X conversion by A. Köhn, using the code IPF-FDMC
- Detailed insight into the mode conversion process

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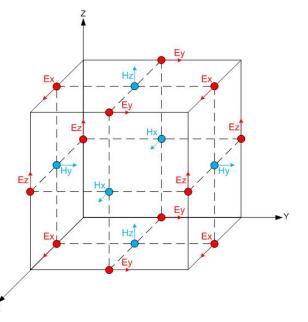
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3D simulations – ADE-FDTD

- New code developed in support of experimental project
- 3D finite difference time-domain (FDTD) method for solving Maxwell's equations :

$$\frac{\partial \mathbf{H}}{\partial t} = -\frac{1}{\mu_0} \nabla \times \mathbf{E} \qquad \qquad \frac{\partial \mathbf{E}}{\partial t} = \frac{1}{\varepsilon_0} \nabla \times \mathbf{H} - \frac{1}{\varepsilon_0} \mathbf{J}$$

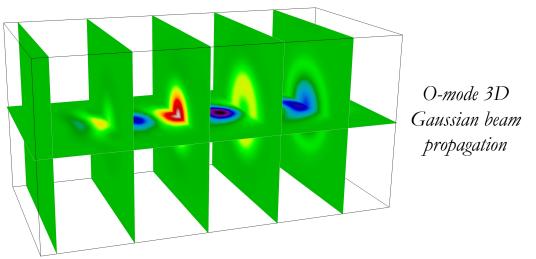
- Discretise field components to staggered grid to simplify calculation of numerical curl
- Substitute in 2nd order centred difference formulae in both space and time
- Obtain leapfrog equations for updating E and B-fields
- For plasma dielectric response, solve: $\frac{\partial}{\partial t} \mathbf{J} = \omega_{pe}^2 \epsilon_0 \mathbf{E} \frac{e}{m} \mathbf{J} \times \mathbf{B}_0$





Code details

- Written in C++. Data-level parallelisation (spatial domain) using MPI
- Arbitrary static n_e and background B_0 profiles specified incident beam then excited using TF/SF source term and simulation run in time domain
- Perfectly matched layer (PML) boundary conditions: very thin absorbing regions
- Full 3D grid output at each timestep very large; virtual sensors reduce output dimensionality / sampling frequency. Transmission coefficients calculated in postanalysis



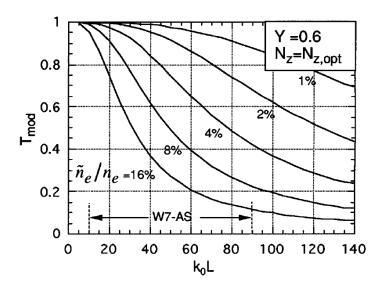
• Future development will include hot plasma terms in order to investigate full B-X-O conversion



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FORM Density fluctuation studies

• Density fluctuations – beam diverges, reduction of conversion efficiency



H.P. Laqua et al, PRL 78 (18), 3467-3470 (1997)

• Analytic modelling - pdf approximating beam divergence used to modify Mjølhus formula

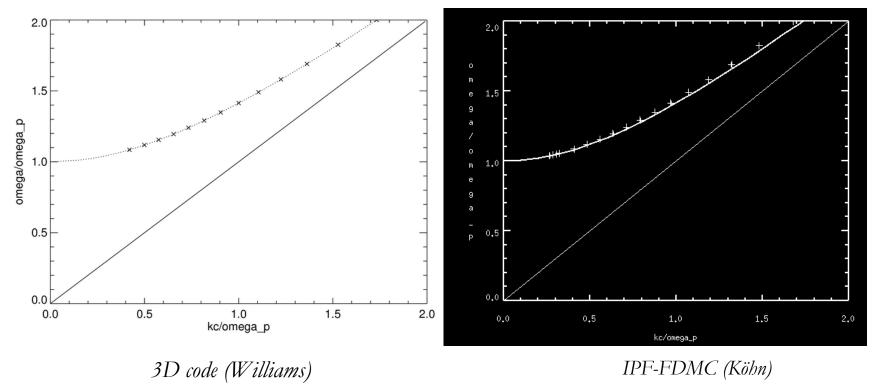
- Initial 1D and 2D full-wave modelling carried out by Köhn using IPF-FDMC
- •3D structures (filaments) at tokamak edge. For oblique incidence, problem is inherently 3D 2D modelling is forced to choose single cut through profile
- Test validity of Laqua result across different regimes

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Code comparison

- New 3-D code compared against results from IPF-FDMC.
- First stage dispersion relations



O-mode dispersion relation : dashed line – analytical, points : numerical

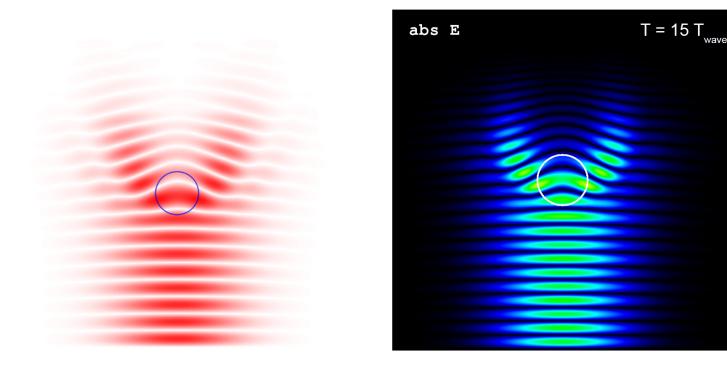
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Code comparison (2)

• 2nd stage – add blob with Gaussian profile, peak density below critical, to homogeneous plasma background



3D code (Williams)

IPF-FDMC (Köhn)

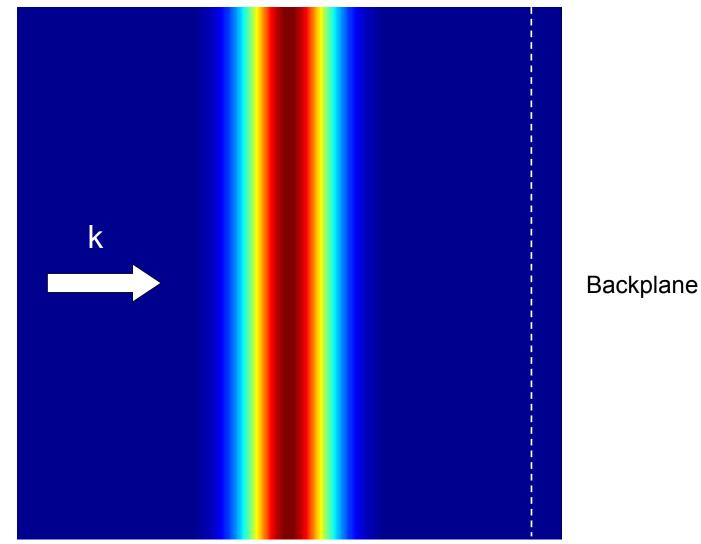
• Circle at location of X = 0.6 surface. Codes agree on beam scattering

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Filament scattering



Incident beam



Analysis

- Average electric field on backplane over several cycles
- Calculate total E, mean and σ in 2 dimensions estimate of 'degree of scattering'

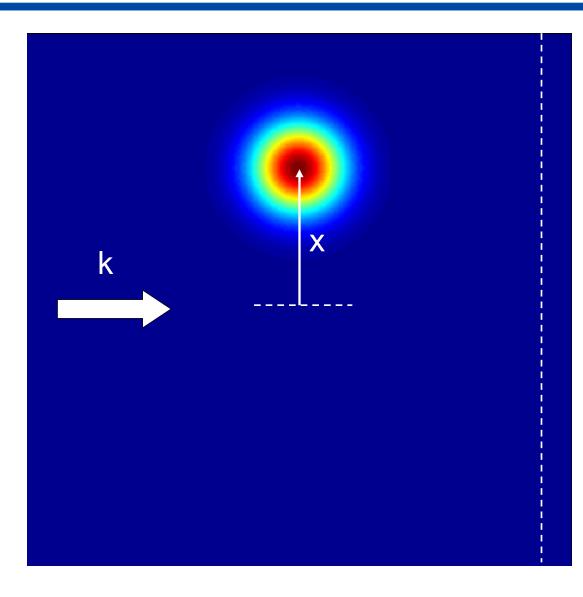
$$\mu = \frac{\int \int \mathbf{x} E_{avg}(x,y) \, dx \, dy}{E_{tot}}$$

$$\sigma = \frac{\int \int (\mathbf{x} - \mu)^2 E_{avg}(x, y) \, dx \, dy}{E_{tot}}$$

• Scan each parameter through experimentally relevant values

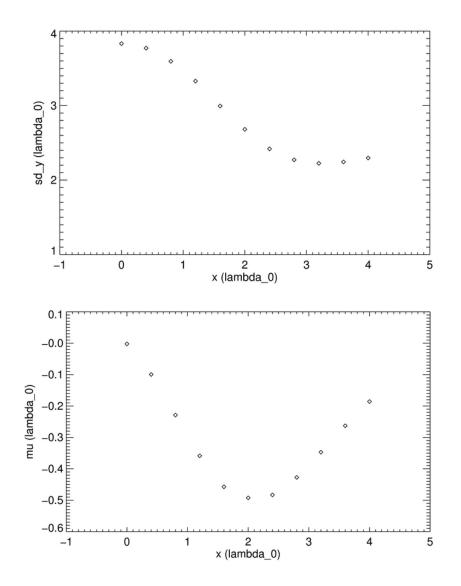


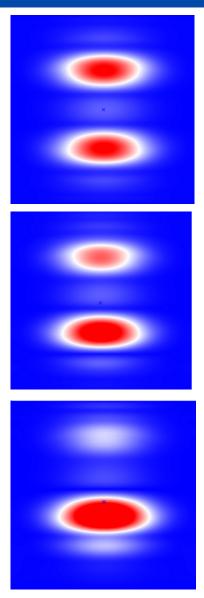
Parameter 1 : position





Position scan results





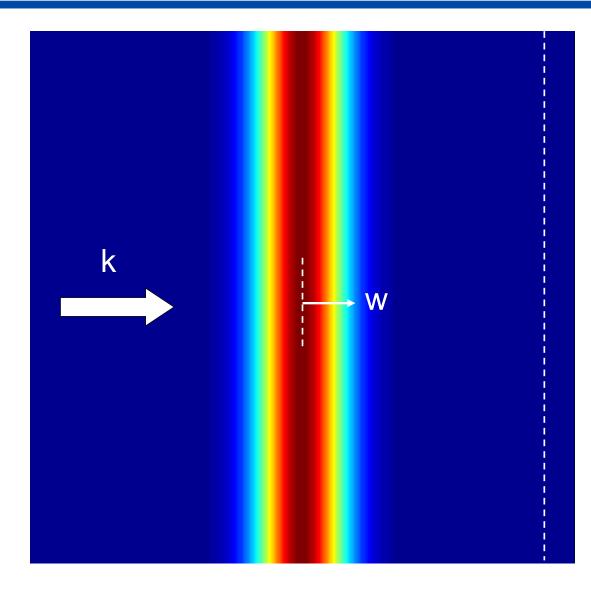
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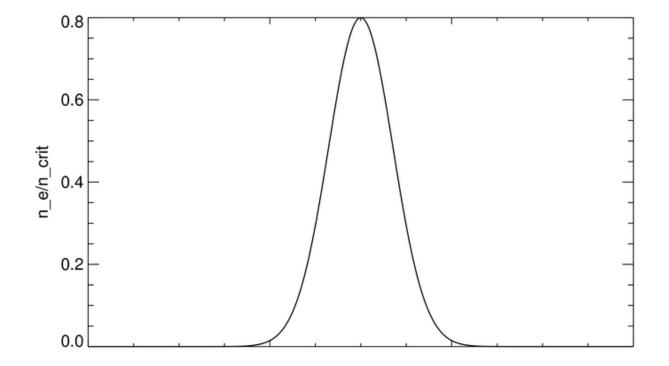


Parameter 2 : width



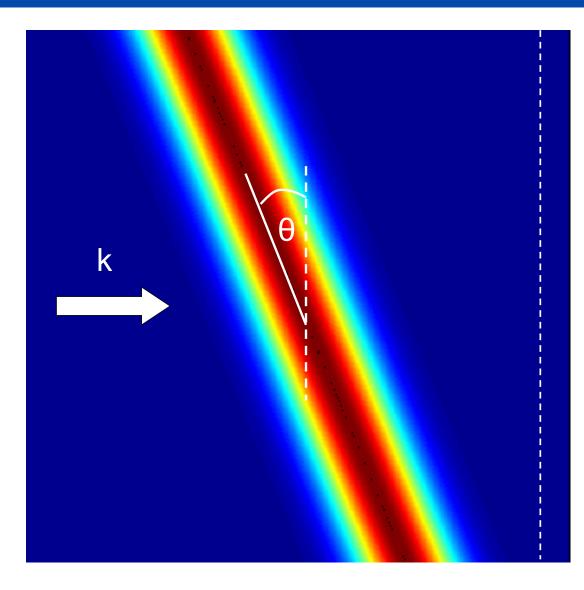


Parameter 3 : density





Parameter 4 : angle





• Investigate the influence of presence of blob near mode conversion surface on mode conversion efficiency – compare against Laqua formula

• Include more realistic turbulent profiles for ST mode conversion region (generated from code e.g. GS2) for highest possible relevance to experimental studies. Average results over a set of perturbed profiles

- Include real MAST experimental profiles for comparison
- Investigate the effect of magnetic shear on mode conversion efficiency

• Can the effect of magnetic fluctuations be distinguished from that of density fluctuations?

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• A new 3D FDTD code has been developed to simulate mode conversion in a fusion plasma.

• This is being used to investigate 3D effects including that of density fluctuations at the turbulent boundary of a spherical tokamak.

• These results are being compared against 2D simulations in a collaboration with IGVP Stuttgart.

• Results will aid the interpretation of data from new MAST diagnostic producing 2D images of mode conversion windows.

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